

Measurement of Methane Autoignition Delays in Carbon Dioxide Diluent at Supercritical Conditions

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1 Introduction and background

The directly fired supercritical carbon dioxide (sCO₂) power cycle enables the dual benefit of high efficiency and nearly complete capture of CO₂. Different from conventional gas turbines, the operating conditions of the sCO₂ power cycle features high CO₂ concentration as diluent and extremely high pressures ranging from 100 bar to 300 bar, above the critical pressure of CO₂. Currently, there are no records on combustion properties at these conditions. This work reports the measurement of autoignition delays of CH₄/O₂/CO₂ mixtures at sCO₂ conditions for the first time and answers several open chemical kinetics questions at sCO₂ conditions: (a) Does chemically reactive CO₂ significantly alter ignition chemistry at high pressure conditions? (b) Are existing kinetic models derived mostly from data obtained at lower pressures capable of predicting autoignition delays at sCO₂ conditions?

In this work, autoignition delays of CH₄/O₂/CO₂ mixtures are measured above the critical point of CO₂ (at an average pressure of 100 bar and in the temperature range of 1274 to 1433 K) in a high-pressure shock tube with large inner diameter (15.24 cm). The large diameter is critical to this study as CO₂ diluent magnifies the non-ideal effect and increases the thickness of the boundary layer. High Mach number required for achieving the above conditions further increases shock attenuation mainly from wall viscous effects [11]. Therefore, a large diameter shock tube is required to avoid interaction between the boundary layer and the center region.

2 Experimental methods

A high pressure shock tube is used to measurement autoignition delays. The shock tube used in this study is fabricated out of 316 stainless steel with 15.24 cm inner diameter and a wall thickness of 5.08 cm. The total length of the shock tube is 20 m consisting of a 10 m long driver section (only 6 m used for this study) and a 10 m long driven section. The inner surface of the tube is electro polished with a 0.2 μm surface finish. The shock tube is equipped with six endwall plugs and eight circumferential ports (located 1.24 cm away

from the endwall) enabling pressure and optical measurements. The maximum operating pressure of the tube is 376 atm with preheating capability up to 500 K. The shock tube is equipped with a hydraulically operated diaphragm ram section and a custom-designed contour valve for quick vacuuming. To damp out vibrations caused by high pressure experiments, the test section is clamped into a 2-ton concrete dead mass anchored to the floor.

The temperature after the reflected shock (T_5) is calculated from normal shock relations using the obtained shock velocity and normally the ideal equation of state (EoS). However, at elevated pressure, the choice of real gas EoS versus ideal gas EoS for the T_5 calculation needs to be examined in detail. Since the pressure after the reflected shock (P_5) is as high as 200 bar, the density prediction using the ideal gas EoS may not represent the actual density after the reflected shock. Davidson *et al.* [9] derived and solved normal shock relations by implementing real gas equations of state: Peng-Robinson, Soave, Redlich-Kwong and Van der Waals for pure argon up to 1000 atm and temperatures between 1000–3000 K. It was found that when using the Peng-Robinson EoS, there was a 83 K/1000 atm difference when compared to an ideal gas. A similar approach is used in this study and showed a minor deviation of ± 4 K in T_5 from ideal gas calculations.

3 Pressure and emission measurements

When diatomic or polyatomic gases are used as diluents, the wall viscous effects (e.g. the thickness of the boundary layer) becomes significantly more noticeable. As a result, the normal portion of the shock wave in the boundary layer turns into a λ -shaped shock near these regions [6], a phenomenon referred to as shock bifurcation. A recent study from Hargis and Petersen [14] showed the significance of the boundary layer effect for gas mixtures with CO_2 , Ar and N_2 dilution. The study reported that the thickness of boundary layer is approximately 30 mm in a N_2/CO_2 mixture at 1400 K and 1.83 bar after 600 μs . Therefore, in similar shock tube experiments, the sidewall pressure trace could be contaminated. Another deleterious consequence of the reflected shock-boundary layer interaction is an observed pressure rise after the passage of the reflected shock (i.e. dP_5/dt). In the current work, dP_5/dt derived from sidewall pressure measurements is between 15–18%/ms as demonstrated in Fig. 1a. Correspondingly, the value from endwall traces is between 3–5%/ms for a 5:10:85 mixture of $\text{CH}_4/\text{O}_2/\text{CO}_2$. During the analysis process, the obtained pressure rise over time is simulated in CHMEKIN-PRO [15] for an unbiased comparison between experimental and numerical results. The pressure rise over time (dP_5/dt) were modeled in CHMEKIN-PRO [15] by applying a time-dependent pressure fit using linear regression of the endwall pressure trace.

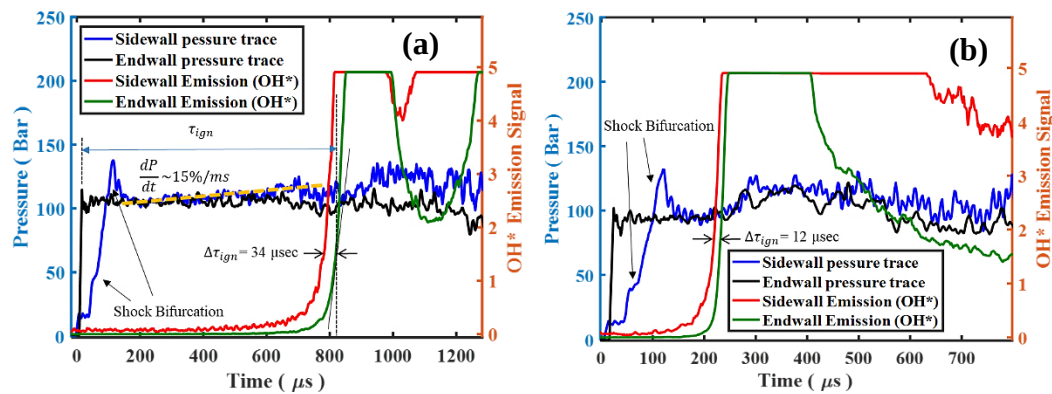


Figure 1. An example of simultaneous sidewall and endwall pressures and OH* emission traces for $\text{CH}_4/\text{O}_2/\text{CO}_2=5:10:85$ at 100 bar and temperature: a) 1274 K and b) 1433 K

To determine the occurrence of autoignition, OH* chemiluminescence is recorded from both sidewall and endwall throughout all experiments. Having fuel concentrations $\geq 2\%$ in the current study, the ignition delay time (τ_{ign}) is defined as the time difference between the arrival of the reflected shock at the endwall and the maximum slope of endwall OH*, as demonstrated in Fig. 1a.

4 Results and discussion

The IDTs for CH₄/O₂/CO₂ mixtures at stoichiometric and a fuel-rich conditions ($\Phi=2$) measured at 100 bar and temperature between 1274–1433 K are shown in Fig. 2. The predicted IDTs from different kinetic models are compared with experiments. The simulation results herein are obtained by solving a zero-dimensional reactor with constant internal energy and volume using CHMEKIN-PRO [15]. While USC Mech II [19], Aramco 2.0 [8, 20-25], HP-Mech [26] and FFCM-1 [27] are able to predict IDTs within the uncertainty of the experimental data, GRI 3.0 [7] underpredicts the autoignition delay by a factor of 3.

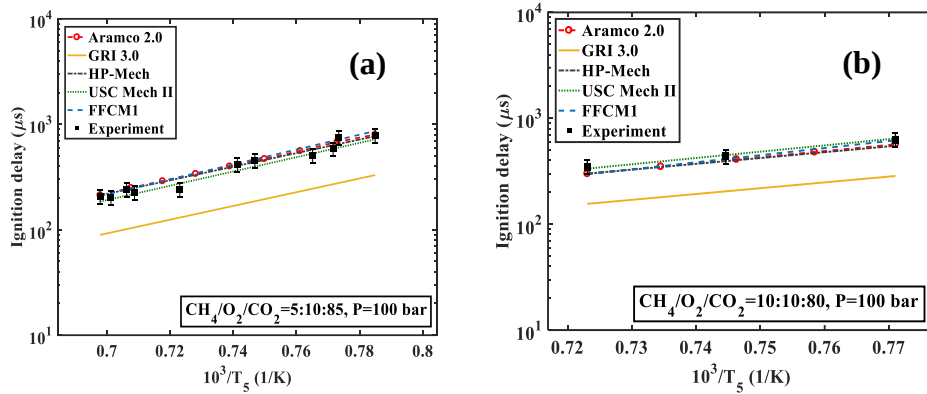


Figure 2. Ignition delay time of CH₄/O₂/CO₂ at 100 bar for (a): stoichiometric ($\Phi=1$), $T=1274$ – 1433 K (b): rich ($\Phi=2$) mixtures, $T=1297$ – 1383 K

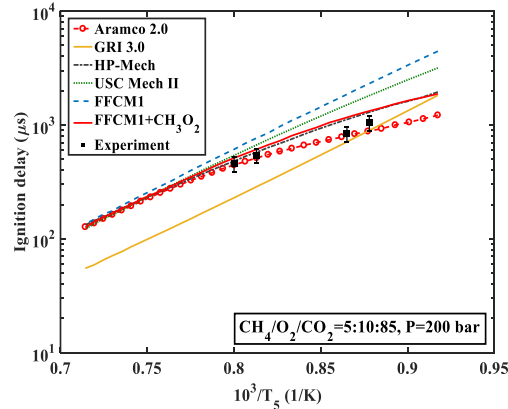


Figure 3. Ignition delay time measurements of stoichiometric CH₄/O₂/CO₂ mixtures at 200 bar, $T=1139$ – 1250 K

Figure 3 summarizes the results obtained from 200 bar experiments for a 5:10:85 mixture of CH₄/O₂/CO₂. At temperatures less than 1250 K, the deviation in IDT predicted by different kinetic models grows while at higher temperature ranges, all kinetic models except for GRI 3.0 [7] show a converging trend. As shown in Fig. 3, the experimental data show very good agreement to IDT predictions from Aramco 2.0 [8, 20-25] for a temperature range of 1139 to 1250 K. HP-Mech [26] also gives a reasonable prediction to the

experimental data. FFCM-1 [27], on the other hand, overpredicts the ignition delay at these conditions by approximately a factor of 2.5. GRI 3.0 [7] still fails to predict the ignition trend at 200 bar, having a maximum deviation of approximately 3 times faster ignition delay.

To better understand the behavior of different kinetic models, a reaction pathway flux analysis is conducted. This analysis produces the ratio of carbon flux through different reaction pathways using a Global Pathway Selection (GPS) method [28] in CANTERA [29]. The results for $\text{CH}_4/\text{O}_2/\text{CO}_2$ (5:10:85) at 1200 K and 200 bar is shown in Fig. 4.

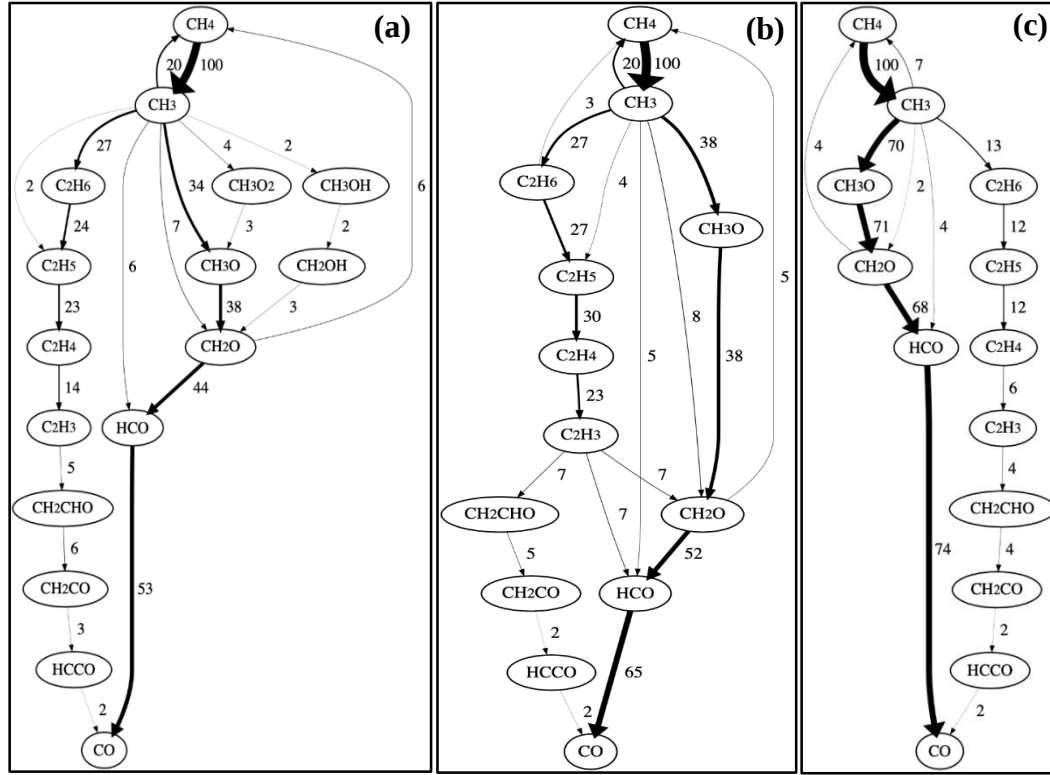
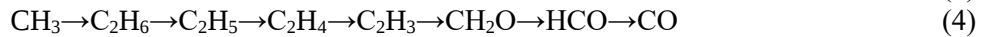
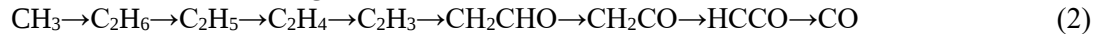


Figure 4. Reaction path flux results for $\text{CH}_4/\text{O}_2/\text{CO}_2=5:10:85$ at 1200 K and 200 bar using (a) Aramco 2.0 (b) FFCM-1 and (c) GRI 3.0

Numbers shown in Fig. 4 represent the percentage of carbon flux (normalized by the total carbon flux of the fuel, CH_4 in this study) through various pathways. After the initiation reaction of CH_4 to form CH_3 (methyl radical), there exist two major reaction pathways of CH_3 . One is CH_3 oxidation which is through:



The second reaction pathway of CH_3 is recombination to form ethane (C_2H_6). Tracking this pathway to formation of CO, pathway (2) for Aramco 2.0 [8, 20-25], pathways (2-4) for FFCM-1 [27] and pathway (2) for GRI 3.0 [7] are observed as following:



The flux ratio between CH_3 oxidation and recombination pathways (i.e. $\text{CH}_3 \rightarrow \text{CH}_3\text{O}$ and $\text{CH}_3 \rightarrow \text{C}_2\text{H}_6$) for Aramco 2.0 [8, 20-25], FFCM-1 [27] and GRI 3.0 [7] are 1.3, 1.4 and 5.4, respectively. The significantly larger CH_3 oxidation flux ratio explains why GRI 3.0 [7] predicts a much faster autoignition than other two kinetics models.

It can also be seen that two new CH_3 reaction pathways are predicted by Aramco 2.0 [8, 20-25]: formation of CH_3O_2 and CH_3OH which could accelerate autoignition. The experimental data for $\text{CH}_4/\text{O}_2/\text{CO}_2$ mixtures at 200 bar shows a unique trend that is only captured by Aramco 2.0 [8, 20-25]. As observed from Fig. 3, this kinetic model shows a different trend for low and high temperatures. The ignition behavior predicted by Aramco 2.0 [8, 20-25] at 200 bar for the $\text{CH}_4/\text{O}_2/\text{CO}_2$ mixture suggests a strong temperature dependency especially at lower temperatures. A distinct difference is observed in reactions that promote ignition at the two different temperatures. One major difference is the reactions of CH_3O_2 with CH_2O and CH_4 , which accelerate oxidation, and therefore promote ignition at 1100 K. At 1400 K, the reaction of CH_3O_2 with CH_3 to form CH_3O becomes relevant for autoignition and the importance of this reaction increases with pressure. Interestingly, the sensitivity analysis performed by Shao *et al.* [10] using FFCM-1 [27] also showed the importance of CH_3O_2 to ignition, while this species is not actually included in FFCM-1 [27]. It can be concluded that the $\text{CH}_3\text{O}_2/\text{CH}_3\text{OH}$ kinetics accounts for the unique ignition behavior predicted by Aramco 2.0 [8, 20-25] at low temperature, 200 bar conditions. To further verify the effect of CH_3O_2 in predicting ignition delays at various temperatures, simulations were performed using an improved version of FFCM-1 with the addition of CH_3O_2 kinetics [30]. The results are also presented in Fig. 3. It can be seen that the improved version of FFCM-1 with $\text{CH}_3\text{O}_2/\text{CH}_3\text{OH}$ addition matches the experiments well. At high pressure and lower temperature conditions, CH_3 recombines with O_2 to form CH_3O_2 , which then oxidizes CH_4 and CH_3 and accelerates ignition.

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