# Study of overdriven detonations of H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixtures

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## 1. Introduction

The specific features of self-sustained detonations of a mixture of fuel with NO<sub>2</sub> as oxidizer is the appearance of two main exothermic reaction steps during the heat release process behind the leading shock wave. These two steps find their expression in the existence of a double detonation cellular structure ( $\lambda_1$  and  $\lambda_2$ ) for the stoichiometric and rich mixtures (and probably for the very lean mixtures). These two structures are correlated to the two calculated characteristic chemical induction lengths (Li<sub>1</sub> and Li<sub>2</sub>) consequence of the existence of two local maximum reaction rates during the heat release behind the leading shock wave (Lamoureux et al., 2001; Joubert, 2001; Joubert et al., 2003; Desbordes et al., 2004). In particular in H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> systems, for equivalence ratio  $\phi \ge 1$  in the range 0.5 <  $\phi$  < 2.0, each two cells  $\lambda_1$  and  $\lambda_2$  of the first and second reaction steps varies proportionally to Li<sub>1</sub> and Li<sub>2</sub> (calculated in the ZND frame with a detailed chemical kinetics (Joubert et al., 2003)) by a quasi constant factor k ( $\lambda_i = k Li_i$ ) of about 15-20. In order to estimate, in these conditions, the post-shock temperature dependency of  $Li_i$  (and  $\lambda_i$ ), overdriven detonation of different overdriven degrees, D/D<sub>CJ</sub>, are considered numerically and experimentally. For these different detonation propagation regimes, the comparison between calculations and experiments of the  $Li_i/Li_{i,CJ}$  (D/D<sub>CJ</sub>) and  $\lambda_i/\lambda_{i,CJ}$  (D/D<sub>CJ</sub>) dependencies can provide global activation energy Ea<sub>i</sub> of each reaction step assuming that Li<sub>i</sub> ~  $\lambda_i$  ~ exp(Ea<sub>i</sub>/RT). Experimentally, quasi-steady overdriven detonation in rich H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixtures are obtained by using the shock tube technique explained in details in Desbordes and Lannoy, 1990.

## 2. Experimental set-up

Detonations of H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixtures were conducted experimentally at ambient initial temperature  $T_0 = 293$ K in a steel detonation tube (52 mm internal diameter) consisting of two parts separated by a thin Mylar foil (thickness of 40µm). The first part A, the driver section, of 6 m long allows initiating and stabilizing a quasi-CJ detonation wave in stoichiometric C<sub>2</sub>H<sub>4</sub>/O<sub>2</sub> mixture at various initial pressure P<sub>0,A</sub>. Provided that P<sub>0,A</sub> is large enough in comparison with P<sub>0,B</sub> (the initial pressure of the part B), the Mylar foil rupture caused by this detonation wave involves the creation of a quasi-steady overdriven detonation wave in the second part (part B), the driven section of 1 m long. This section contains the H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixture at equivalence ratio  $\phi = 1.2$  and initial pressures P<sub>0,B</sub> = 0.4 bar and 0.5 bar.

Instrumentation includes pressure transducers in order to characterize the detonation propagation regimes (detonation pressure jump P and velocity D) in the first and the second part of the tube and smoked foil plates in the driven section in order to record cellular

structure of the overdriven detonation wave in H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixtures. For fixed initial pressure P<sub>0,B</sub> of H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub>, initial pressure P<sub>0,A</sub> of C<sub>2</sub>H<sub>4</sub> + 3O<sub>2</sub> mixtures was determined with respect to the expected overdriven degree D/D<sub>CJ</sub> of the detonation regime in H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixtures given by the detonation polar analysis in the P<sub>1</sub>, U<sub>1</sub> plane on Fig. 1 (see Desbordes and Lannoy, 1990).



#### 3. Results

Steady overdriven detonations are obtained in H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixtures quasi instantaneously by choosing adequate condition of initial pressure P<sub>0,A</sub> in the driver section. Typical smoked foil records of cellular structure in H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixture at  $\phi = 1.2$ , P<sub>0,B</sub> = 0.4 bar are displayed in Fig. 2 for the CJ detonation regime (Fig. 2a) and two steady overdriven detonation regimes characterized by their overdriven degree (D/D<sub>CJ</sub> ratio) D/D<sub>CJ</sub> = 1.07 (Fig. 2b) and D/D<sub>CJ</sub> = 1.12 (Fig. 2c).

The two nets of cellular structure of size  $\lambda_1$  (small) and  $\lambda_2$  (large) appear clearly. The  $\lambda_1$  cell widths are deduced from these records and are equal to 1.4 mm, 0.32 mm and 0.18 mm for  $D/D_{CJ} = 1.00$ ,  $D/D_{CJ} = 1.07$  and  $D/D_{CJ} = 1.12$  respectively while the  $\lambda_2$  cell widths are found to be equal to 42.9 mm ( $D/D_{CJ} = 1.00$ ), 13.7 mm ( $D/D_{CJ} = 1.07$ ) and 10.2 mm ( $D/D_{CJ} = 1.11$ ).



**Figure 2.**: Typical smoked foil records for CJ detonation regime (a) and for 2 steady overdriven detonation conditions:  $D/D_{CJ} = 1.07$  (b) and  $D/D_{CJ} = 1.12$  (c) in H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixture at  $\phi = 1.2$ , P<sub>0,B</sub> = 0.4 bar and T<sub>0</sub> = 293K.

As previously observed in single detonation cellular system (Vasiliev and Nikolaev, 1978; Gravrilenko and Prokhorov, 1983; Huang and Xu, 1987; Desbordes, 1988; Desbordes et al., 1990; Meltzer et al., 1990), the cell sizes  $\lambda_1$  and  $\lambda_2$  drastically decreased as D/D<sub>CJ</sub> ratio increased. As D/D<sub>CJ</sub> increased from 1 to about 1.1, cellular structure sizes decrease from 1 to roughly 0.3. The overdriven cell sizes have decreased drastically up to a factor around of 10 in comparison of those measured in the quasi-CJ detonation regime.

Experimental values  $\lambda_i / \lambda_{i,CJ}$  vs. D/D<sub>CJ</sub> and numerical values Li<sub>i</sub>/Li<sub>i,CJ</sub> vs. D/D<sub>CJ</sub> for P<sub>0,B</sub> = 0.4 bar and P<sub>0,B</sub> = 0.5 bar in H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixtures are reported for each reaction step ( $\lambda_1$ , Li<sub>1</sub> and  $\lambda_2$ , Li<sub>2</sub>) in Fig. 3 and Fig. 4 respectively. Chemical induction lengths (Li<sub>1</sub> or Li<sub>2</sub>) are numerically derived from the calculated induction times of the detailed chemical kinetic mechanism (Joubert et al., 2003) computed with the STANJAN code of the CHEMKIN II package (Kee et al., 1993; 1994).



Agreements between numerical  $\text{Li}_i/\text{Li}_{i,\text{CJ}}$  (D/D<sub>CJ</sub>) and experimental  $\lambda_i/\lambda_{i,\text{CJ}}$  (D/D<sub>CJ</sub>) curves are quite good with only a small difference between these results essentially due to the measure

uncertainty on detonation cellular widths. In a general way, the size  $\lambda_i$  of the detonation structure of a chemical system depends on the local detonation velocity D by the approximate relationship:

$$\frac{\lambda_i}{\lambda_{i,CJ}} \cong \frac{Li_i}{Li_{i,CJ}} \cong \frac{D}{D_{CJ}} \exp\left(\frac{Ea_i}{RT_{ZND}} \left[ \left(\frac{D_{CJ}}{D}\right)^2 - 1 \right] \right)$$

where  $\lambda_i = \lambda_{i,CJ}$  when  $D = D_{CJ}$ , Ea<sub>i</sub> being the activation energy of the step (i).

The slope of the curve  $\lambda_i/\lambda_{i,CJ} = f(D/D_{CJ})$  deduced of this expression is approximately equal to  $-2Ea_i/RT$ . In the particular case  $D/D_{CJ} = 1.0$ , the slope of the curves  $\lambda_1/\lambda_{1,CJ} = f(D/D_{CJ})$  corresponds to  $-2Ea_1/RT_{ZND}$  and that of  $\lambda_2/\lambda_{2,CJ} = f(D/D_{CJ})$  corresponds to  $-2Ea_2/RT_1$ ,  $T_1$  being the temperature at the end of the first exothermic step and at the beginning of the second exothermic step. So, it is possible to deduce the reduced activation energy (Ea<sub>i</sub>/RT) and the activation energy (Ea<sub>i</sub>) from the two curves, each of them can be associated respectively with the two characteristic chemical steps during the process of heat release. Those results are presented in Table 1.

$P_{0,B} = 0.4$ bar	Slope of $Li_i/Li_{i,CJ}$ (D/D <sub>CJ</sub> ) curves	Т	Ea <sub>i</sub> /RT	Ea <sub>i</sub> (kcal/mol)
Fig 3a	-12.74	$T_{ZND} \approx 1725 K$	$\approx 6.37$	$\approx 21.8 (i = 1)$
Fig 3b	-13.50	$T_1 \approx 2938 K$	$\approx 6.75$	$\approx$ 39.4 (i = 2)
$P_{0,B} = 0.5$ bar	Slope of $Li_i/Li_{i,CJ}$ (D/D <sub>CJ</sub> ) curves	Т	Ea <sub>i</sub> /RT	Ea <sub>i</sub> (kcal/mol)
Fig 4a	-12.74	$T_{ZND} \approx 1733 K$	$\approx 6.37$	$\approx 21.9 (i = 1)$
Fig 4b	-13.31	$T_1 \approx 2946 K$	≈ 6.66	$\approx 39.0$ (i = 2)

**Table 1.:** Activation energies (Ea<sub>i</sub>) of the two exothermic reaction steps for H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixture deduced from the results of the Li<sub>i</sub>/Li<sub>i,CJ</sub> =  $f(D/D_{CJ})$  curves at  $\phi = 1.2$ , P<sub>0,B</sub> = 0.4 bar and 0.5 bar, T<sub>0</sub> = 293K



The activation energy (respectively Ea<sub>1</sub> ~ 21.8 kcal/mol at P<sub>0,B</sub> = 0.4 bar and Ea<sub>1</sub> ~ 21.9 kcal/mol at P<sub>0,B</sub> = 0.5 bar) of the first exothermic step (H<sub>2</sub>+NO<sub>2</sub>  $\rightarrow$  NO+H<sub>2</sub>O) presents a value close to the computed one for a H<sub>2</sub>/O<sub>2</sub> mixture, *i.e.* Ea ~ 19 kcal/mol according to the study of Desbordes et Lannoy (1990) or a value close to the numerical one deduced from the Li<sub>i</sub>/Li<sub>i,CJ</sub> =  $f(D/D_{CJ})$  curves on Fig. 5 (with Ea ~ 15 kcal/mol) for a H<sub>2</sub>/O<sub>2</sub> mixture in the same conditions. The activation energy (respectively Ea<sub>1</sub> ~ 39.4 kcal/mol at P<sub>0,B</sub> = 0.4 bar and Ea<sub>2</sub> ~ 39.0 kcal/mol at P<sub>0,B</sub> = 0.5 bar) of the second exothermic step (H<sub>2</sub>+NO  $\rightarrow$   $\frac{1}{2}N_2$ +H<sub>2</sub>O) presents a value close to the computed one for a H<sub>2</sub>/NO mixture (Ea ~ 34.5 kcal/mol deduced from Li<sub>i</sub>/Li<sub>i,CJ</sub> =  $f(D/D_{CJ})$  curves on Fig. 6) at  $\phi = 1.2$  and T<sub>0</sub> = 293K.

## 4. Conclusion

The study of different overdriven detonation regimes in H<sub>2</sub>-NO<sub>2</sub>/N<sub>2</sub>O<sub>4</sub> mixture of equivalence ratio  $\phi = 1.2$  has revealed two cellular structures of different size  $\lambda_1$  and  $\lambda_2$  which decreases drastically when overdriven degree (D/D<sub>CJ</sub>) increases. A quite good agreement between numerical Li<sub>i</sub>/Li<sub>i,CJ</sub> (D/D<sub>CJ</sub>) and experimental  $\lambda_i/\lambda_{i,CJ}$  (D/D<sub>CJ</sub>) curves is observed for both cellular structures. The deduced activation energy of the first exothermic step (H<sub>2</sub>+NO<sub>2</sub>  $\rightarrow$ NO+H<sub>2</sub>O), Ea<sub>1</sub> ~ 21.8 kcal/mol, is close to the value of the activation energy (Ea ~ 15 - 19 kcal/mol) of H<sub>2</sub>/O<sub>2</sub> mixture and the activation energy of the second exothermic step (H<sub>2</sub>+NO  $\rightarrow \frac{1}{2}N_2$ +H<sub>2</sub>O), Ea<sub>2</sub> ~ 39.0 kcal/mol, is close to the activation energy for H<sub>2</sub>/NO mixture (Ea ~ 34.5 kcal/mol). This study confirms the existence for overdriven detonation regime of two detonation cell structures and two exothermic reaction steps that we had already identified in quasi-CJ detonation regime.

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