## A weak solution of point source blast wave problem

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The "point source model" consists of a system of gasdynamics equations that describes the motion of ideal gas initiated by an instantaneous release of finite amount of energy at a point (or line or plane) in a uniform state. This model has originally been introduced to describe the blast wave propagation caused by an intense explosion<sup>1)</sup>. The original system of equations has conveniently been converted into a system of equations for four unknowns  $f(x, y), g(x, y), h(x, y), \lambda(y)$  on the region  $0 \le x, y \le 1$  by the use of blast wave transformation<sup>1)</sup>.

Here we introduce the variables  $\xi = -\log x$  and  $\eta = \log R$ , R being related to  $\lambda$  by  $\lambda = (R/y)dy/dR$ . With  $T = (xR)^{\alpha+1}/y$ , we also set F = f/x,  $G = gh^{-\gamma+1}T$ , H = hyT,  $I = (\gamma hf^2/2 + g/(\gamma - 1))T$ , where  $\alpha = 0, 1$  and 2 correspond respectively to plain, line, and point sources, and  $\gamma$  is the ratio of the specific heats of the gas. We convert the of equations for  $f, g, h, \lambda^{(1)}$  into the form

$$-U_{\xi} + V_{\eta} = 0, \ U(0,\eta) = -R^{\alpha+1}K, \quad 0 \le \xi < \infty, \ -\infty < \eta < \infty,$$
(1)

where

$$U = \begin{pmatrix} (F-1)I + FG(\frac{H}{Ty})^{\gamma-1} \\ (F-1)G \\ (F-1)H \end{pmatrix}, V = \begin{pmatrix} I \\ G \\ H \end{pmatrix}, K = \begin{pmatrix} \frac{1}{\gamma-1} \\ (1 + \frac{\gamma-1}{\gamma+1}(1-y))(1 - \frac{2}{\gamma+1}(1-y))^{\gamma}y^{-1} \\ 1 \end{pmatrix}.$$
 (2)

This system of equations is supplemented by the condition (conservation of energy)

$$\int_0^1 \frac{I}{x} dx = R_0^{\alpha+1} + \frac{1}{(\alpha+1)(\gamma-1)} R^{\alpha+1},$$

where  $R_0$  is the characteristic length of the explosion defined by the explosion energy and the initial pressure.

We consider a weak solution of (1), that is, (U, V) satisfying

$$\int_{-\infty}^{\infty} U_i(0,\eta) \Phi(0,\eta) d\eta + \int_0^{\infty} \int_{-\infty}^{\infty} (U_i \Phi_{\xi} - V_i \Phi_{\eta}) d\eta d\xi = 0 \quad (i = 1, 2, 3)$$
(3)

for all smooth functions  $\Phi$  which are identically zero outside some bounded set<sup>2</sup>). If  $\eta = \varphi(\xi, s)$  is a solution, depending on *i*, of the ordinary differential equation  $d\eta/d\xi = V_i/U_i$  with  $\eta|_{\xi=0} = s$  for U, V defined in (2), then changing the variables  $\xi, \eta$  in (3) to s, t determined by  $\xi = t, \eta = \varphi(\xi, s)$  shows that  $U_i = U_i(0, s)/\varphi_s$  and  $V_i = U_i\varphi_{\xi}$  satisfy (3). In practice, we may construct a solution from an explicit expression for  $V_i/U_i$  in variables  $\xi$  and  $\eta$ .

## References

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